

# Energy Correlators for Jet Analysis

Richard Lin

Relativistic Heavy Ion Group (RHIG)



## Introduction to the n-point projected energy correlator

$$\mathcal{E}(\vec{n}) = \lim_{r \rightarrow \infty} \int_0^{\infty} dt \, r^2 n^i T_{0i}(t, r\vec{n}), \quad (1)$$

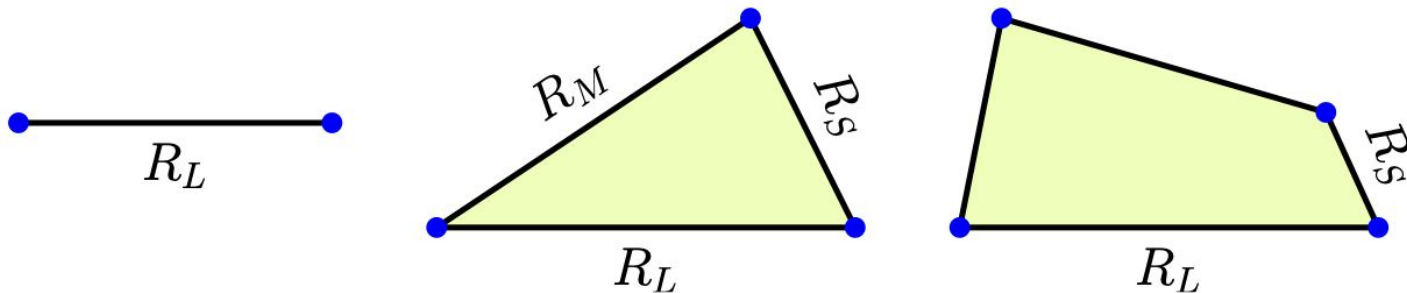
[\(Komiske, Moult, Thaler, Zhu, 2023\)](#)

$$\begin{aligned} \text{ENC}(R_L) &= \left( \prod_{k=1}^N \int d\Omega_{\vec{n}_k} \right) \delta(R_L - \Delta \hat{R}_L) \\ &\cdot \frac{1}{(E_{\text{jet}})^N} \langle \mathcal{E}(\vec{n}_1) \mathcal{E}(\vec{n}_2) \dots \mathcal{E}(\vec{n}_N) \rangle, \end{aligned} \quad (2)$$

$$\text{ENC}(R_L) = \sum_{i_1, \dots, i_N} \int dR_L \frac{p_T^{i_1} p_T^{i_2} \dots p_T^{i_N}}{p_{T, \text{jet}}^N} \delta(R_L - \Delta \hat{R}_L)$$

## Introduction to the n-point projected energy correlator

N=2 and N=3:



[\(Komiske, Moul, Thaler, Zhu, 2023\)](#)

## Introduction to the n-point projected energy correlator

### Looking ahead:

- another way to probe  $\alpha_s$
- in the regions where perturbative theory is reliable:

$$\frac{E3C(R_L)}{EEC(R_L)} \sim \alpha_s + \mathcal{O}(\alpha_s^2) + \dots$$

- a new way to measure a fundamental constant from jet observables

# Introduction to the n-point projected energy correlator

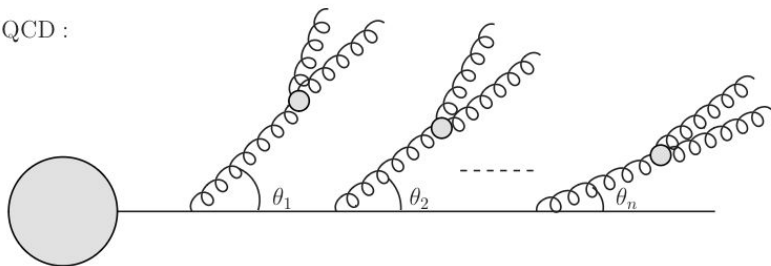
## IRC safety

1. Collinear and infrared (soft gluon) emissions are physically unresolvable

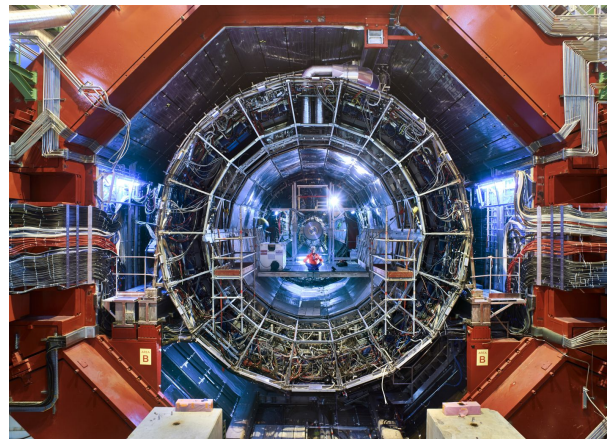
$N \equiv$  number of jet constituents

$$z_{\text{leading}} \equiv \frac{p_{T, \text{leading}}}{\sum_{i \in \text{jet}} p_{T,i}}$$

QCD :



[Image Credits: Luizoni and Marzani, 2015.](#)



## Introduction to the n-point projected energy correlator

### The pest of dimensionality

1. The standard n-point energy correlator function (ECF):

$$e_n^{(\beta)} = \sum_{i_1 < i_2 < \dots < i_n} \left( \prod_{a=1}^n p_{T, i_a} \right) \prod_{1 \leq a < b \leq n} R_{i_a i_b}^{\beta}$$

2. one dimensional only for the n=2 case. At n=3 need to compare three unique values for R -> 3-axes needed for visualization. At n=4, would need a 6 dimensional plot!
3. generated data also get sparser, which can complicate analysis: the curse of dimensionality

# Methods

## Jet Matching

1. A key complication in detector studies
2. A theorist's definition of jets involves perturbative QCD calculations including all particles
3. Detectors measure the tracks of charged particles, but neutral particles can only be seen by calorimeters -> jets are constructed using charged particles
4. Also have to account for detector effects: hadronization, pileup, noise, etc.

$$\begin{aligned}
 & \sigma \rightarrow \rho \quad \left[ -g_{\mu\nu} + (1-\xi) \frac{p_\mu p_\nu}{p^2 + i0} \right] \frac{i}{p^2 + i0} \\
 & \quad \frac{i(p + m_f)_{\rho\sigma}}{p^2 - m_f^2 + i0} \\
 & \quad \frac{i}{p^2 + i0} \\
 & b, \sigma \rightarrow a, \rho \quad -ig\mu^\epsilon (t^\alpha)_{ab} \gamma_\rho^\mu \\
 & \quad \mu, \alpha \\
 & \alpha \rightarrow q, \gamma \quad -g\mu^\epsilon f_{\alpha\beta\gamma} q^\mu \\
 & \quad \beta, \mu \\
 & p, \alpha, \lambda \rightarrow q, \beta, \mu \quad -g\mu^\epsilon f_{\alpha\beta\gamma} \left[ (p-q)^\nu g^{\lambda\nu} + (q-r)^\lambda g^{\mu\nu} + (r-p)^\mu g^{\nu\lambda} \right] \\
 & \quad r, \gamma, \nu \\
 & \alpha, \kappa \rightarrow \beta, \lambda \quad \begin{aligned} & -ig^2 \mu^{2\epsilon} f_{\alpha\beta\gamma} f_{\gamma\delta\epsilon} (g^{\kappa\mu} g^{\lambda\nu} - g^{\kappa\nu} g^{\lambda\mu}) \\ & -ig^2 \mu^{2\epsilon} f_{\alpha\gamma\delta} f_{\delta\beta\epsilon} (g^{\kappa\lambda} g^{\mu\nu} - g^{\kappa\nu} g^{\lambda\mu}) \\ & -ig^2 \mu^{2\epsilon} f_{\alpha\delta\beta} f_{\beta\gamma\epsilon} (g^{\kappa\lambda} g^{\mu\nu} - g^{\kappa\nu} g^{\lambda\mu}) \end{aligned} \\
 & \quad \delta, \nu \rightarrow \gamma, \mu
 \end{aligned}$$

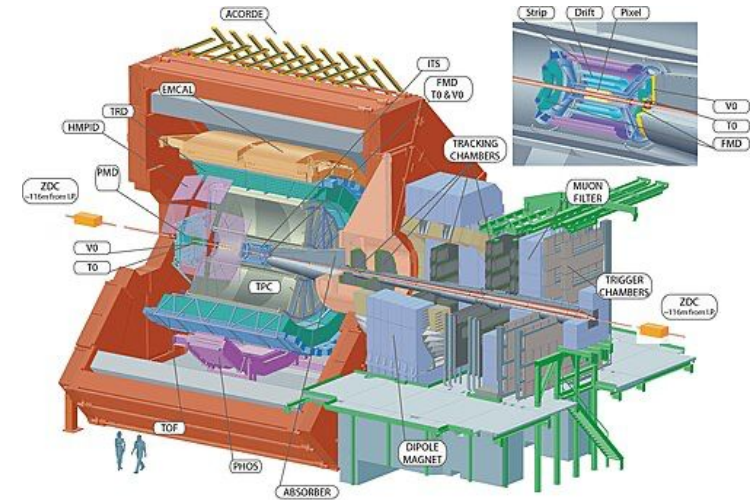
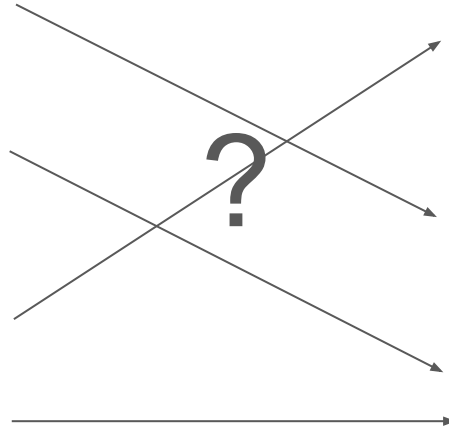


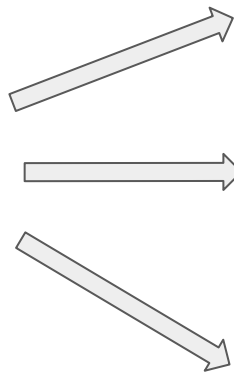
Image Credits: ALICE collaboration

Methods

Simulation pipeline



PYTHIA event  
generation

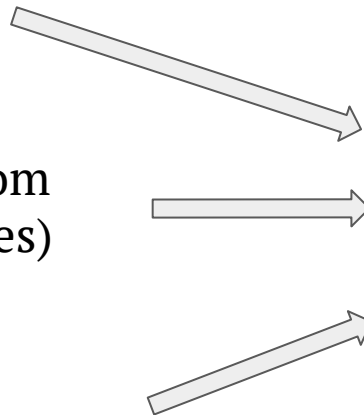


# FastJet

Full jet

Charged Jet (from  
charged particles)

Charged Jet  
reconstructed



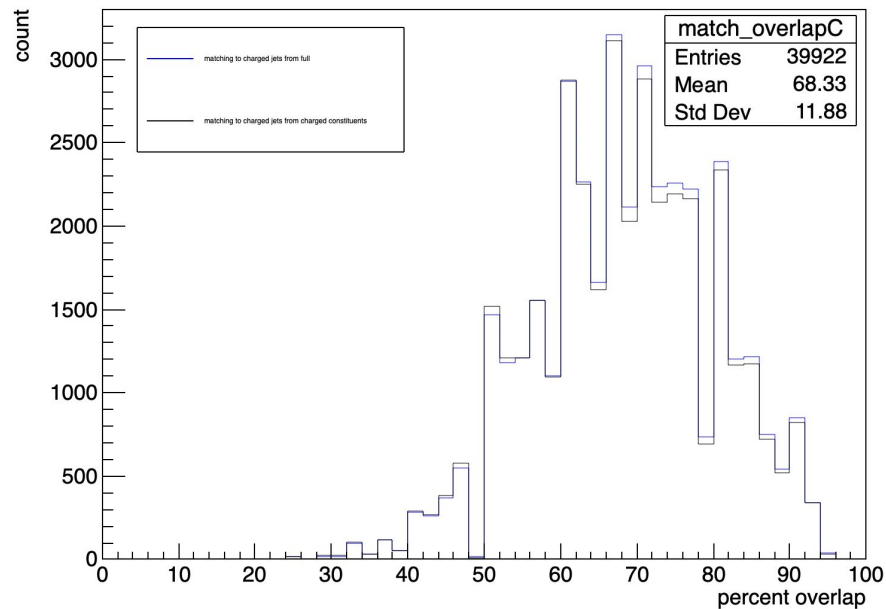
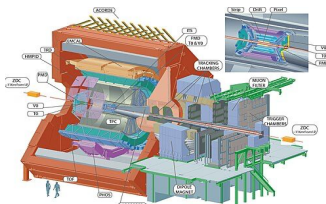
Analysis:  
matching,  
correlators, ...



## Methods

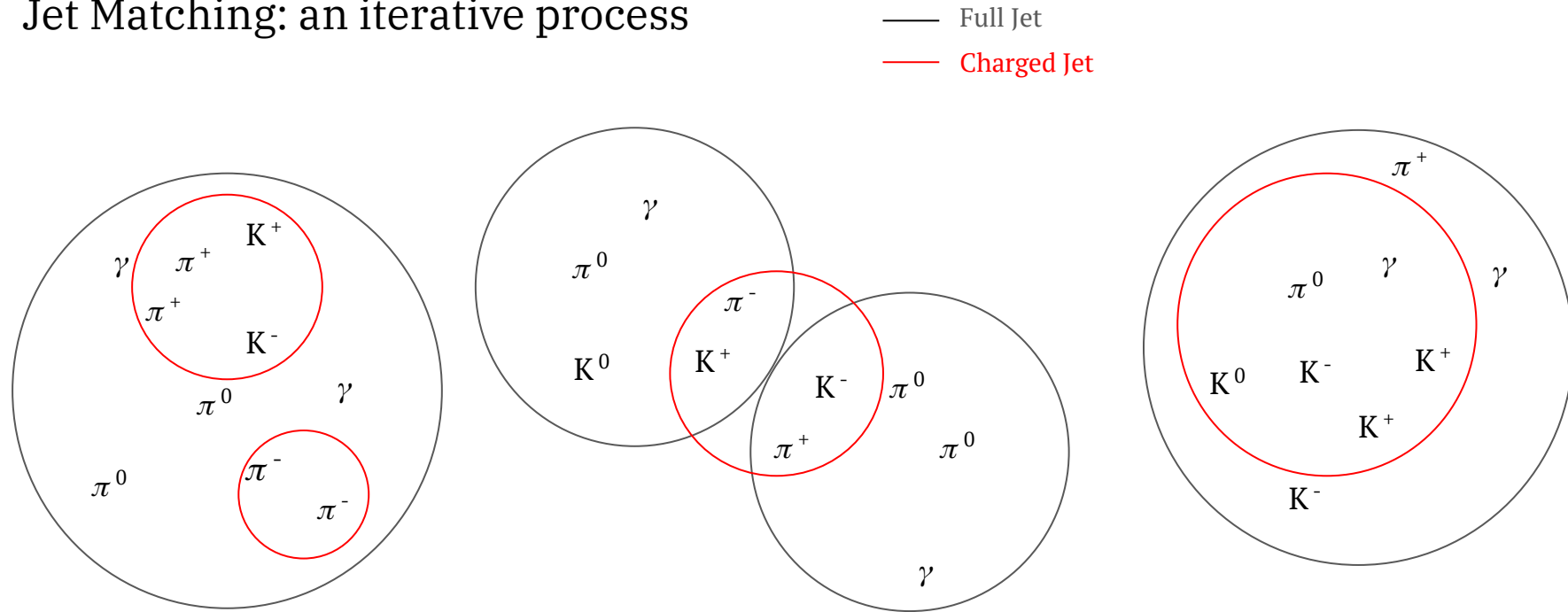
## Jet Matching: an iterative process

[illegible]



## Methods

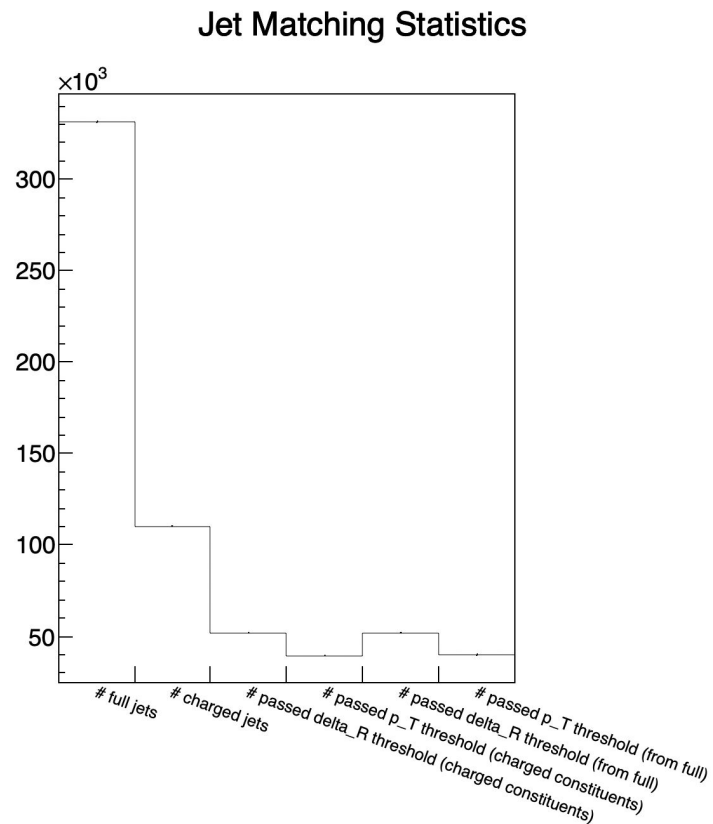
### Jet Matching: an iterative process



1. Cut first based on difference in jet axis ( $R < 0.1$ )
2. Further cut based on minimum constituent pT fraction of charged jet compared to full ( $pT_{\text{charged}}/pT_{\text{Full}} > 0.8$ )

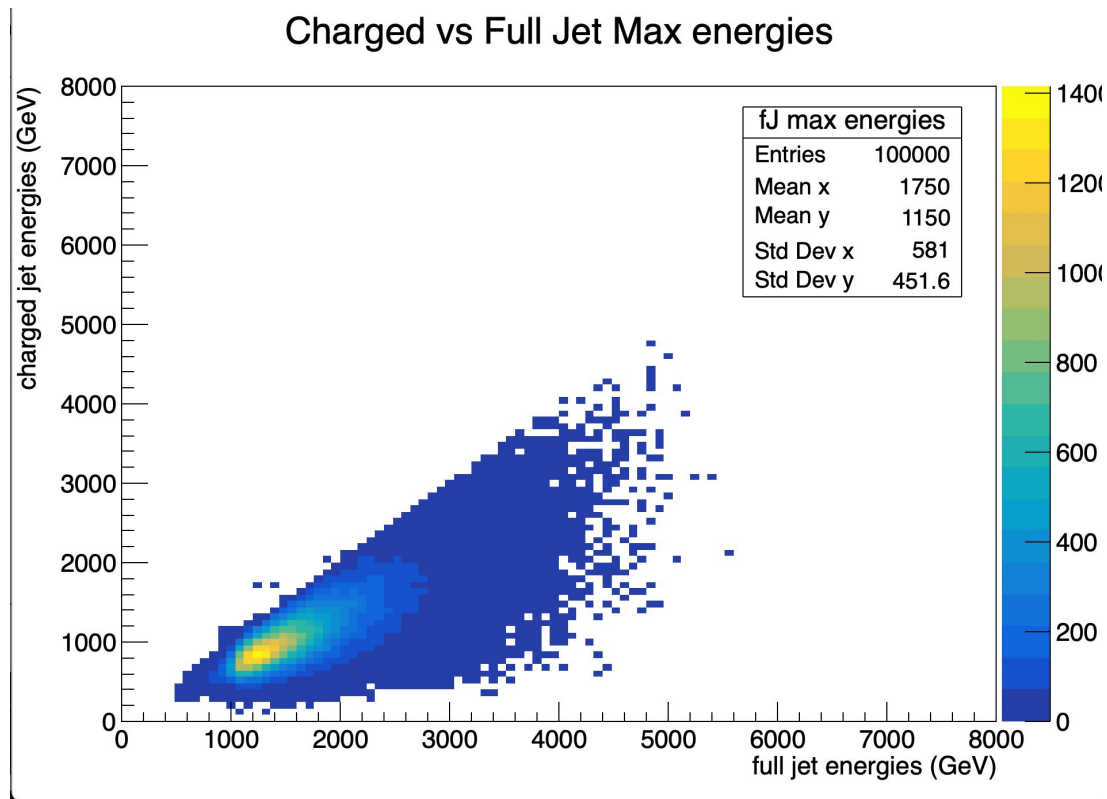
## Methods

# Jet Matching: an iterative process



## Methods

### Jet Matching: an iterative process



# Results

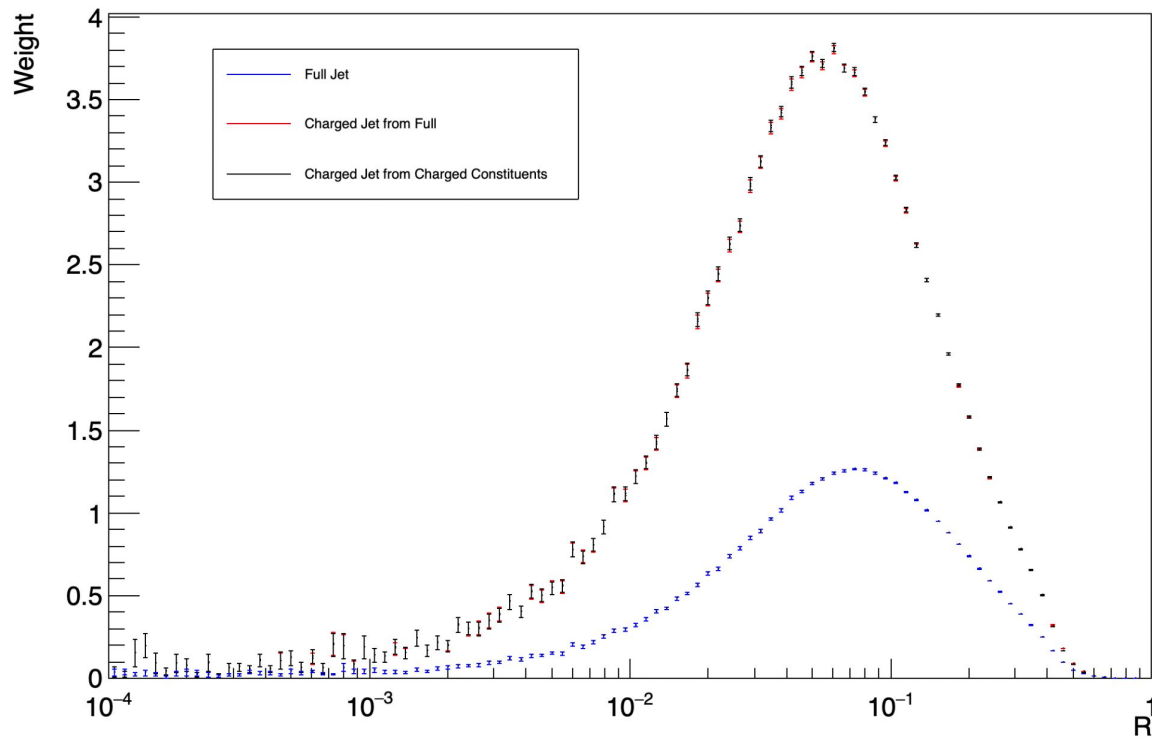
## Energy correlators from Pythia

$$\sqrt{s} = 13.0 \text{ TeV}$$

$$R = 0.4, \eta_{jet} < 0.5$$

$$p_{T,track} > 1 \text{ GeV}/c$$

EEC (N=2)



## Results

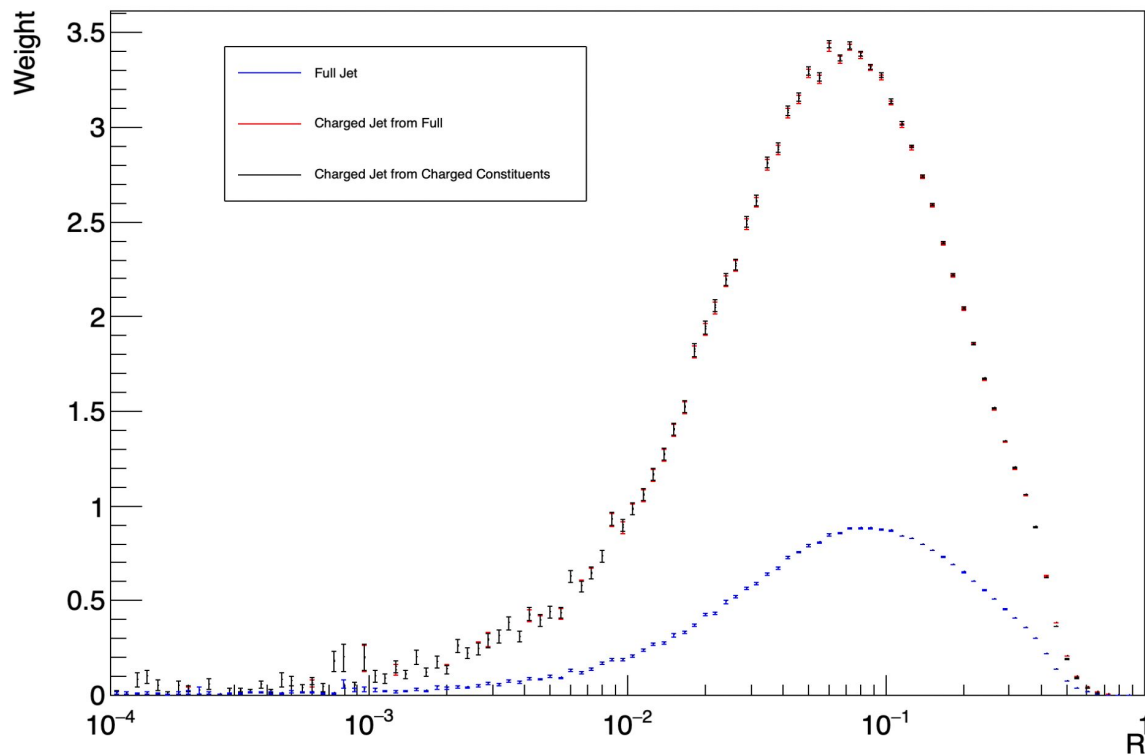
### Energy correlators from Pythia

$$\sqrt{s} = 13.0 \text{ TeV}$$

$$R = 0.4, \eta_{jet} < 0.5$$

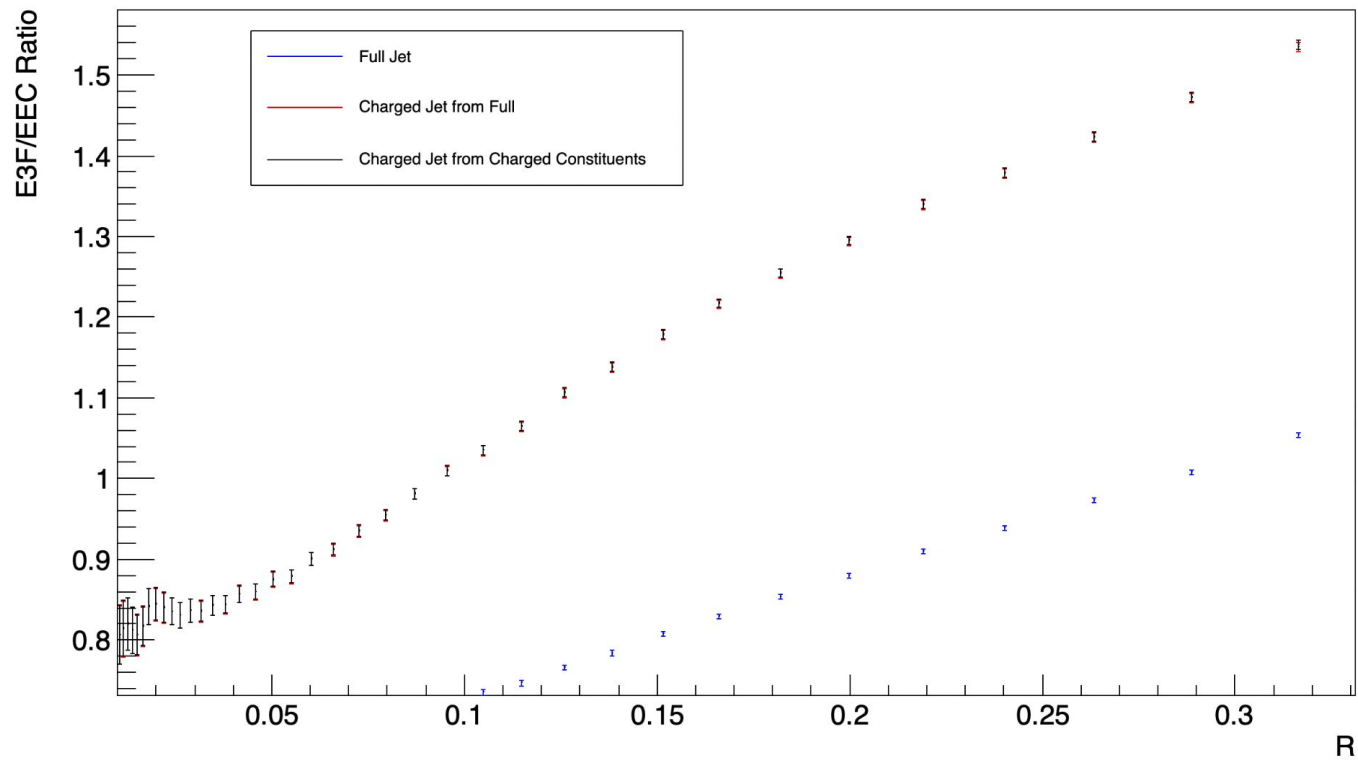
$$p_{T,track} > 1 \text{ GeV}/c$$

E3C (N=3)



## Results

### Energy correlators from Pythia



## Discussion

### Future directions

- n-point energy correlators ( $n > 3$ )
- extracting  $\alpha_s$
- tuning jet matching
- beyond pp collisions